

## Introduction

- Standard quantum key distribution (QKD) is limited to about 250 km due to losses in the optical fiber.
- Quantum repeaters [Bri1998] permit to extend this distance by nested entanglement distillation and entanglement swapping.
- The secret key rate (bits per memory per second) resulting from a quantum repeater is given by

$$K = \frac{R r_\infty}{M}, \quad (1)$$

where

- $R$  (repeater rate) is the average number of generated entangled pairs per second,
- $r_\infty$  is the secret fraction, i.e., the ratio of the secret bits and the measured bits in the asymptotic limit (Devetak-Winter bound  $1 - S(X|E) - H(X|Y)$ ),
- $M$  is half the number of used memories per repeater node.
- We investigate the quantum repeater with encoding [Jia2009] in the context of quantum key distribution and compare it to the quantum repeater using distillation, as the former does not require classical communication.

## Generic quantum repeater

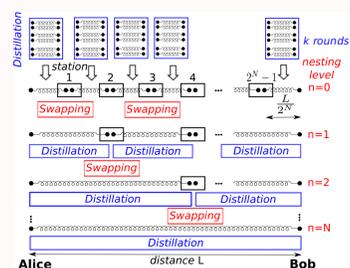


Fig. 1: A generic quantum repeater protocol [Bri1998] with maximal nesting level  $N$  and  $k$  rounds of distillation in all nesting levels.

- Problem: classical communication is needed for acknowledging the success of entanglement distribution, distillation and swapping.

## Quantum repeater with encoding

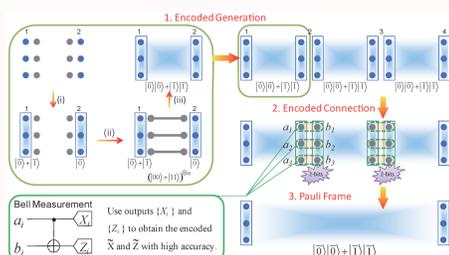


Fig. 2: Repeater protocol with encoding, from [Jia2009].

- Advantage: classical communication is only needed for acknowledging the success of entanglement distribution and in the end for communicating the Pauli frame.
- Disadvantage: many logical gates are needed.

## QR with encoding: remote CNOT and error models

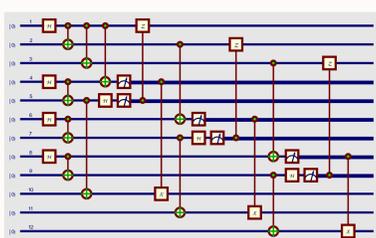


Fig. 3: Remote CNOT for the quantum repeater with encoding, adapted from [Jia2007].

- Application of multiple two-qubit gates and neglecting errors of order  $\beta^2 = (1 - p_G)^2$  and higher leads to

$$(1 - \text{Length}[op](1 - p_G)) \bigotimes_{j=1}^{\text{Length}[op]} op[j] \rho \left( \bigotimes_{j=1}^{\text{Length}[op]} op[j] \right)^\dagger + (1 - p_G) \left\{ \sum_{i=1}^{\text{Length}[op]} \bigotimes_{j=1}^{i-1} op[j] f \left( i, \rho, \bigotimes_{j=i+1}^{\text{Length}[op]} op[j] \right) \left( \bigotimes_{j=1}^{i-1} op[j] \right)^\dagger \right\},$$

where  $op = \{U_m, \dots, U_1\}$  is the list of gates and  $f(i, \rho, A) := \text{tr}_i(A \rho A^\dagger) \otimes \frac{1}{4}$ .

Assumptions:

- One-qubit operations are error free,
- error model for two-qubit operations (depolarizing map):  $O^{real} \rho = p_G O^{ideal} \rho + \frac{1-p_G}{4} \mathbb{1}$ ,
- Bell pairs are depolarized:

$$\rho_{Dep}(F_0) := F_0 \Pi_{|\phi^+\rangle} + \frac{1-F_0}{3} (\Pi_{|\phi^-\rangle} + \Pi_{|\psi^+\rangle} + \Pi_{|\psi^-\rangle}).$$

## The repeater rate

- Average number of attempts to connect  $m$  pairs, each generated with probability  $P_0$  ( $P_0 = 10^{-\alpha L_0/10}$  is the probability that a photon is not absorbed at a distance  $L_0 = L/m$ ) and deterministic entanglement swapping [Ber2011]:

$$Z_m(P_0) := \sum_{j=1}^m \binom{m}{j} \frac{(-1)^{j+1}}{1 - (1 - P_0)^j}. \quad (2)$$

### Generic Quantum Repeater

- The repeater rate including the classical communication time can be found in [Bra2013].
- Using distillation and no classical communication time the rate is [Abr2013]:

$$R_{Rep} = \frac{1}{2T_0} \left( \frac{2}{3} \right)^{N+\sum_n k_n} P_0 \prod_{n=1}^N P_{ES}(n) \prod_{i=0}^{k_n} P_D^O(i, n), \quad (3)$$

$P_D^O(i, n)$  is the success probability in the  $i$ -th distillation round and  $n$ -th nesting level for the *Oxford* protocol [Deu1996],  $T_0 = L_0/c$  ( $c$  is the speed of light in the optical fiber), and  $P_{ES}(n)$  is the success probability of entanglement swapping in the  $n$ -th nesting level.

### Quantum repeater with encoding

- For deterministic swapping:

$$R_{QEC} = \frac{1}{T_0 Z_{nm}(P_0)}, \quad (4)$$

with  $n$  being number of physical qubits to encode one logical qubit.

## Memories

### Quantum repeater with distillation:

- Number of needed memories depend on the distillation protocol:

- recursive protocol (*Oxford protocol* [Deu1996]):  $M_O = 2 \sum_i k_i$ ,
- entanglement pumping (*Innsbruck protocol* [Due1999]):  $M_I = N + 2 - |\{k_i : k_i = 0\}|$ .

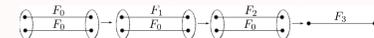


Fig. 4: Entanglement pumping (*Dür et al.* protocol [Due1999]) with  $k = 3$  rounds of purification.

- For optimality of the distillation protocols and strategies see [Bra2013].

### Quantum repeater with encoding

- Number of memories used here is  $M_{enc} = 2n$  (overhead for the remote CNOT).

## Results: optimal repeater protocol

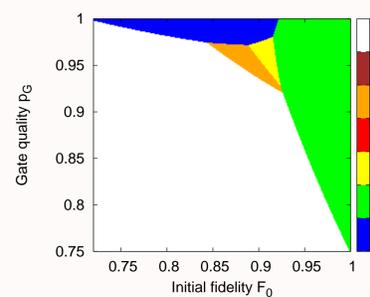


Fig. 5: Optimal quantum repeater protocols w.r.t. the secret key rate per memory per second for  $N = 1$  in terms of the initial fidelity  $F_0$  and the gate quality  $p_G$ .

- Here: distillation only in the end ( $\vec{k} = \{0, k\}$ ) with protocols *Oxford* (O) and *Innsbruck* (I); quantum repeater with encoding (QEC) for the three-qubit repetition code ( $n = 3$ ).
- The number of generated Bell pairs is kept the same, in case of the QEC it is 3, for distillation either 2 (for  $k = 1$ ) or 4 (for  $k = 2$ ).
- The total distance is  $L = 600$  km.
- For initial fidelities  $F_0 \leq 0.85$  the QEC is optimal.
- For an initial fidelity above  $F_0 = 0.92$ , no distillation is optimal.
- The *Innsbruck protocol* is not optimal for this set of parameters, but it was shown in [Bra2013] that this can be achieved for other parameters.

## Results: optimal secret key rate

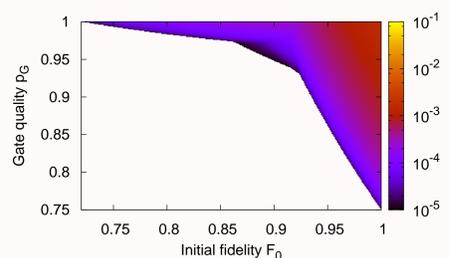


Fig. 6: Optimal secret key rate per memory per second for the quantum repeater protocols ( $N = 1$ ) shown above in terms of the initial fidelity  $F_0$  and the gate quality  $p_G$ .

- The generation of a non-zero secret key rate for  $N = 1$  at a distance of  $L = 600$  km is limited to very good gates and good initial fidelities.
- For having good gates ( $p_G \geq 0.98$ ), but modest fidelities ( $F_0 \leq 0.8$ ), we can still obtain a secret key rate per memory per second on the order of  $10^{-4}$ .

## Discussion

- We calculated the secret key rate per memory per second by comparing two approaches for the quantum repeater: either using distillation or using quantum error correction.
- We found that for modest fidelities ( $F_0 \leq 0.8$ ) we can still obtain a non-zero secret key rate, but we require good gates ( $p_G \geq 0.98$ ).
- Future work includes the extension of these calculations to higher nesting levels (more swappings) and other error correcting codes.

## Acknowledgments

We thank S. Abruzzo, M. Epping, and L. Jiang for fruitful discussions. The authors acknowledge financial support by the German Federal Ministry of Education and Research (BMBF, project QuOReP).

## References

- [Abr2013] S. Abruzzo, S. Bratzik *et al.*, Phys. Rev. A **87**, 052315 (2013).
- [Ber2011] N. K. Bernardes, L. Praxmeyer, and P. van Loock, Phys. Rev. A **83**, 012323 (2011).
- [Bra2013] S. Bratzik, S. Abruzzo, H. Kampermann, and D. Bruß, Phys. Rev. A **87**, 062335 (2013).
- [Bri1998] H. J. Briegel, W. Dür, J. I. Cirac, and P. Zoller, Phys. Rev. Lett. **81**, 5932 (1998).
- [Deu1996] D. Deutsch *et al.*, Phys. Rev. Lett. **77**, 2818 (1996).
- [Due1999] W. Dür, H. J. Briegel, J. I. Cirac, and P. Zoller, Phys. Rev. A **59**, 169 (1999).
- [Jia2007] L. Jiang, J. Taylor, A. Sørensen, and M. D. Lukin, Phys. Rev. A **76**, 062323 (2007).
- [Jia2009] L. Jiang *et al.*, Phys. Rev. A **79**, 032325 (2009).