# A solid-state quantum interface between stationary and flying qubits

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<u>Thanks to</u>:

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# Spin-photon quantum interface

- <u>Goal</u>: to use single-photon pulses to link (distant) quantum nodes. Applications:
  - quantum repeaters
  - distributed quantum information processing
- <u>Resource</u>: indistinguishable photonic qubits (= the same spatio-temporal profile, center frequency & polarization) or entangled spin-photon pairs

$$|\psi\rangle = (\uparrow, H\rangle + |\downarrow, V\rangle)/\sqrt{2}$$

H,V could denote any «internal» degree of freedom (color, polarization, orbital angular momentum, etc) of the photon

# Outline

- A bright source of indistinguishable single photons
- Creation of quantum entanglement between a single photon and a condensed matter spin
- Teleportation from a propagating qubit to a solid-state spin

## Solid-state spins & emitters

- Solid-state emitters (artificial atoms) can be used to realize high brightness long-lived single-photon sources:
  - no need for trapping
  - easy integration into a directional (fiber-coupled) cavity
  - up to 10<sup>9</sup> photons/sec with >70% efficiency

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  - easy integration into a directional (fiber-coupled) cavity
  - up to 10<sup>9</sup> photons/sec with >70% efficiency
- Three different type of emitters:
  - rare-earth atoms embedded in a solid matrix (Er in glass)
  - Deep defects (NV centers in diamond)
  - Shallow defects in semiconductors (quantum dots)

<u>Note</u>: While the concepts & techniques apply to a wide range of solid-state emitters, we focus on quantum dots

## Quantum dots

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- Neutral quantum dots (QD) are ideal for generation of single and entangled indistinguishable photons, thanks to near-transform limited emission lines.
- Single-electron charged QDs allow for realization of a quantum interface between electron spin and generated photon via spin-state dependent light scattering, leading to spin-photon entanglement.

## Quantum dot Spectroscopy



#### How do we make sure that a light pulse contains a single photon: Photon correlations from a single QD

- Intensity (photon) correlation function:  $g^{(2)}(\tau) = \frac{\langle I(t)I(t+\tau) \rangle}{\langle I(t) \rangle^2}$
- To measure g<sup>(2)</sup>(τ), photons from a quantum emitter are sent to a Hanbury-Brown Twiss setup



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- To measure g<sup>(2)</sup>(τ), photons from a quantum emitter are sent to a Hanbury-Brown Twiss setup
- Photon correlations from a weak pulsed laser (<n> ~ 1); detection of a photon does not change the likelihood of detecting a second.



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- To measure g<sup>(2)</sup>(τ), photons from a quantum emitter are sent to a Hanbury-Brown Twiss setup
- Single quantum emitter driven by a pulsed laser: absence of a center peak indicates that none of the pulses have > 1 photon (Robert, LPN).
- $\Rightarrow$  Signature of a single-photon source



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• Two completely indistinguishable single-photon pulses incident on a beam-splitter never lead to coincidences at the output due to a quantum interference effect.



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- The single photon pulses have to have the same spatio-temporal profile, center frequecy, polarization.
- Indistinguishability ensures the absence of entanglement of single photons with uncontrolled degrees of freedom.

#### A single-photon frequency-qubit from a QD: $|\psi\rangle = a|b|ue\rangle + \beta|red\rangle$

In a neutral QD, the elementary optical excitations are excitons (X0); the two linearly polarized exciton X0 lines are split due to electron-hole exchange by ~ 5 GHz

![](_page_14_Figure_2.jpeg)

By controlling the pulse-shape, detuning and polarization of the resonant laser, we could generate a single-color photon or a two-color photonic qubit

# Interference of photonic qubits (superposition of blue and red photons) coming from two quantum dots

![](_page_15_Figure_1.jpeg)

#### Quantum dots and spin qubits: Faraday geometry ( B<sub>ext</sub> = B<sub>z</sub>)

![](_page_16_Figure_1.jpeg)

 $\Gamma$ : spontaneous emission rate

 $\Omega$ : laser coupling (Rabi) frequency

QD with a spin-up (down) electron only absorbs and emits σ+ (σ-) photons – a recycling transition similar to that used in trapped ions.
 ⇒ Measurement of a spin qubit: |ψ> = α|↑> + β|↓>

# Single-shot measurement of electron spin

![](_page_17_Figure_1.jpeg)

- Prepare the electron spin in
  |↑> or |↓>
- Apply a 0.8 µs resonant laser pulse on the trion transition corresponding to |↓>
- Single-shot measurement fidelity ~ 80% in 0.8 μs
- Fidelity is limited by spin pumping into | ↑> - long duration of excitation leads to initialization of the qubit.

Optical transition from a quantum dot spin qubit in Voigt geometry ( $B_{ext} = B_x$ )

Excitation of a trion state results in either emission of a H polarized red photon to  $|\downarrow\rangle$  state or a V polarized blue photon to  $|\uparrow\rangle$  state, with equal probability.

![](_page_18_Figure_2.jpeg)

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![](_page_19_Figure_2.jpeg)

Similar results by Yamamoto, Steel groups; earlier work by Monroe, Lukin

## Procedure for spin-photon entanglement generation

![](_page_20_Figure_1.jpeg)

![](_page_21_Figure_1.jpeg)

An additional  $\pi$ -pulse (dashed curve) is applied to realize a heralded measurement in the spin-up state.

![](_page_22_Figure_1.jpeg)

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Identical (unconditional) counts for red and blue photons confirm the selection rules.

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The g(2) measurement shows that for the [1.2ns, 1.64ns] time range, probability of two-photon emission is negligible.

![](_page_24_Figure_1.jpeg)

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The g(2) measurement shows that for the [1.2ns, 1.64ns] time range, probability of two-photon emission is negligible.

A spin down (up) measurement event ensures that the detected photon is red (blue).

F1=0.87 ± 0.05 in the computational basis measure.ment

## Measurement of quantum correlations

![](_page_25_Figure_1.jpeg)

- An additional  $\pi/2$  or  $3\pi/2$ pulse (dashed curve) is applied to measure the spin in  $|\uparrow\rangle \pm |\downarrow\rangle$ .

## Measurement of quantum correlations

-

![](_page_26_Figure_1.jpeg)

- An additional  $\pi/2$  or  $3\pi/2$ pulse (dashed curve) is applied to measure the spin in  $|\uparrow\rangle \pm |\downarrow\rangle$ >.
  - The data in b & c shows the coincidence measurement when  $\pi/2$ -pulse is applied.

$$|\tilde{\Phi}\rangle = \frac{1}{\sqrt{2}} (|\omega_{red}\rangle e^{-i\omega_z(t_1 - t_g)} - i|\omega_{blue}\rangle)$$

⇒ Coherent oscillations in conditional detection demonstrate quantum correlations between spin and photon

## Measurement of quantum correlations

![](_page_27_Figure_1.jpeg)

- An additional  $\pi/2$  or  $3\pi/2$ pulse (dashed curve) is applied to measure the spin in  $|\uparrow > \pm |\downarrow >$ .
- The data in b & c shows the coincidence measurement when  $\pi/2$ -pulse is applied.
- The data in d & e shows the coincidence measurement when 3 π/2-pulse is applied.
- F2=0.46  $\pm$  0.04 in the rotated basis measurement; overall fidelity F = 0.67  $\pm$  0.05

W. Gao et al. Nature (2012)

## Teleportation from a photonic qubit to a solid-state spin qubit

• Using spin-photon entanglement as a resource, we can transfer the quantum state of a flying photon onto a confined spin (W. Gao, Nat. Comm. (2013)

![](_page_28_Figure_2.jpeg)

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![](_page_29_Figure_2.jpeg)

<u>Next step</u>: probabilistic entanglement of two distant spins

## Entanglement of distant spins

• We need spins with long coherence time: hole spin

![](_page_30_Figure_2.jpeg)

## Entanglement of distant spins

- Erasing which-path information in single-photon scattering from distant spins, leads to entanglement upon detection.
- Proposal by Cabrillo et al. Phys. Rev. A 59, 1025 (1999)

![](_page_31_Figure_3.jpeg)

## Future: Integrated spin photonics

**Spin-photon entanglement** 

**On-chip manipulation and detection** 

![](_page_32_Picture_3.jpeg)

![](_page_32_Picture_4.jpeg)

## Outlook

- Spin-photon quantum interface with decoherencefree spin qubits (singlet-triplet states in QDs)
- Demonstration of nearly deterministic source of entangled photons using neutral QDs